

MEE5114 (Sp22) Advanced Control for Robotics

Lecture 2: Rigid Body Configuration and Velocity

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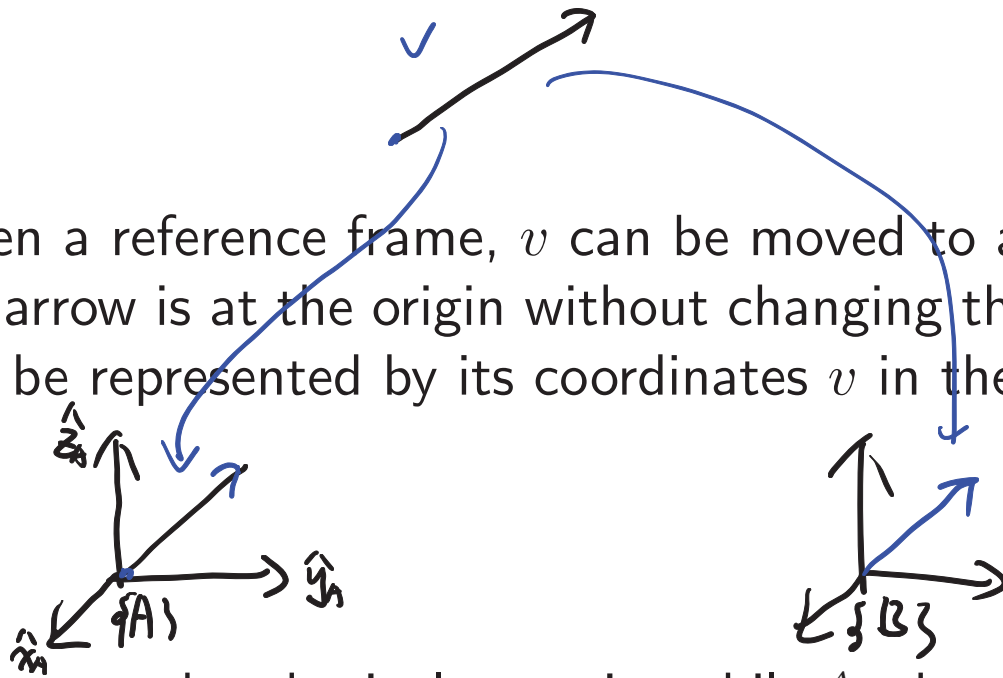
Outline

- Rigid Body Configuration ←
- Rigid Body Velocity (Twist) €
- Geometric Aspect of Twist: Screw Motion

Free Vector

- **Free Vector**: geometric quantity with length and direction

- Given a reference frame, v can be moved to a position such that the base of the arrow is at the origin without changing the orientation. Then the vector v can be represented by its coordinates v in the reference frame.



- v denotes the physical quantity while ${}^A v$ denote its coordinate wrt frame $\{A\}$.

Frame: coordinate sys based on basis vectors

$\{A\}$ - frame: $\{\hat{x}_A, \hat{y}_A, \hat{z}_A\}$.

$${}^A \hat{x}_A = \begin{bmatrix} 1 \\ 0 \\ 0 \end{bmatrix}$$

$${}^A v = \begin{bmatrix} 1 \\ 2 \\ 3 \end{bmatrix} \text{ means}$$

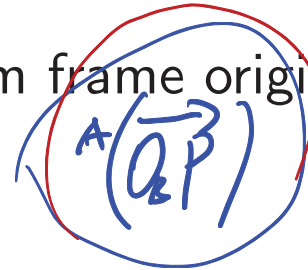
$$v = 1 \cdot x_A + 2 \cdot y_A + 3 \cdot z_A$$

Point

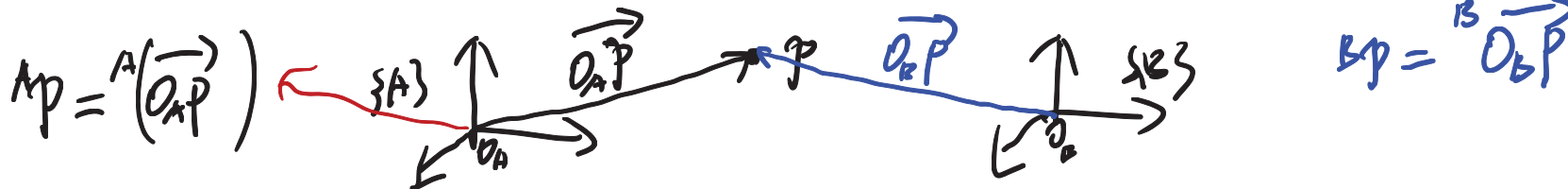
- **Point:** \underline{p} denotes a point in the physical space

• p

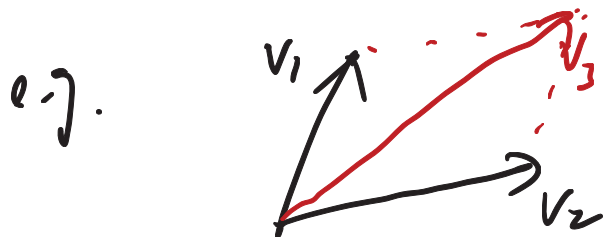
- A point p can be represented by a vector from frame origin to p



- ${}^A p$ denotes the coordinate of a point p wrt frame $\{A\}$



- When left-superscript is not present, it means the physical vector itself or the coordinate of the vector for which the reference frame is clear from the context. think in "coordinate-free" way whenever possible



- coordinate-free : $v_3 = v_1 + v_2$
- But also express "physics" in different frames
- ${}^A v_3 = {}^A v_1 + {}^A v_2$; ${}^B v_3 = {}^B v_1 + {}^B v_2$

Cross Product

$${}^A V_3 \neq {}^B V_1 + {}^B V_2$$

- **Cross product** or **vector product** of $a \in \mathbb{R}^3, b \in \mathbb{R}^3$ is defined as

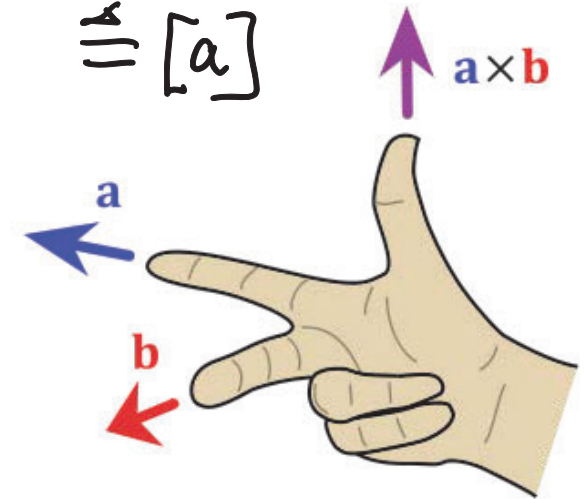
$$a \times b = \begin{bmatrix} a_2 b_3 - a_3 b_2 \\ a_3 b_1 - a_1 b_3 \\ a_1 b_2 - a_2 b_1 \end{bmatrix} = \begin{bmatrix} 0 & -a_3 & a_2 \\ a_3 & 0 & -a_1 \\ -a_2 & a_1 & 0 \end{bmatrix} \begin{bmatrix} b_1 \\ b_2 \\ b_3 \end{bmatrix} \quad (1)$$

$a = \begin{bmatrix} a_1 \\ a_2 \\ a_3 \end{bmatrix}, b = \begin{bmatrix} b_1 \\ b_2 \\ b_3 \end{bmatrix}$

$\stackrel{\Delta}{=} [a]$

Properties:

- $\|a \times b\| = \|a\| \|b\| \sin(\theta)$
- $a \times b = -b \times a$
- $a \times a = 0$



Skew symmetric representation

$$A = A^T$$

- It can be directly verified from definition that $a \times b = [a]b$, where

$$[a] \triangleq \begin{bmatrix} 0 & -a_3 & a_2 \\ a_3 & 0 & -a_1 \\ -a_2 & a_1 & 0 \end{bmatrix} \quad (2)$$

- $a = \begin{bmatrix} a_1 \\ a_2 \\ a_3 \end{bmatrix} \leftrightarrow [a]$

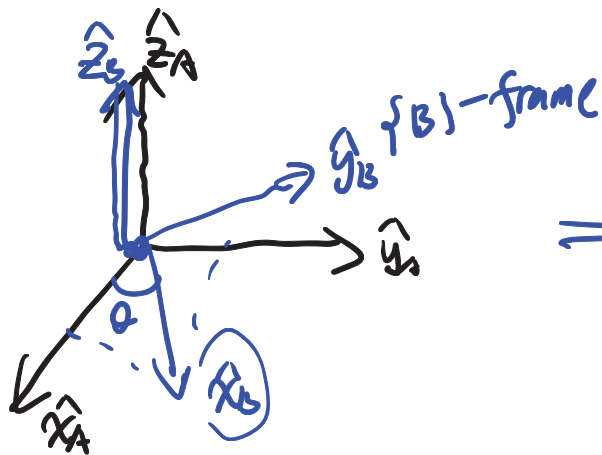
$$a = \begin{bmatrix} 1 \\ 2 \\ 3 \end{bmatrix}, \quad [a] = \begin{bmatrix} 0 & -3 & 2 \\ 3 & 0 & -1 \\ -2 & 1 & 0 \end{bmatrix}$$

- $[a] = -[a]^T$ (called skew symmetric) $\Leftrightarrow A = -A^T$
- $[a][b] - [b][a] = [a \times b]$ (Jacobi's identity)

Rotation Matrix

- **Frame:** 3 coordinate vectors (unit length) $\hat{x}, \hat{y}, \hat{z}$, and an origin
 - $\hat{x}, \hat{y}, \hat{z}$ mutually orthogonal
 - $\hat{x} \times \hat{y} = \hat{z} \Leftarrow$ (right hand rule)
- **Rotation Matrix:** specifies orientation of one frame relative to another

$${}^A R_B \stackrel{\circ}{=} \begin{bmatrix} {}^A \hat{x}_B & {}^A \hat{y}_B & {}^A \hat{z}_B \end{bmatrix}$$



$$\Rightarrow {}^A R_B = \begin{bmatrix} \cos\theta & -\sin\theta & 0 \\ \sin\theta & \cos\theta & 0 \\ 0 & 0 & 1 \end{bmatrix}$$

- A valid rotation matrix R satisfies: (i) $R^T R = I$; (ii) $\det(R) = 1$

$$\begin{bmatrix} {}^A \hat{x}_B^T \\ {}^A \hat{y}_B^T \\ {}^A \hat{z}_B^T \end{bmatrix} \begin{bmatrix} {}^A \hat{x}_B & {}^A \hat{y}_B & {}^A \hat{z}_B \end{bmatrix} = \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix} \quad \left| \begin{array}{l} \det({}^A R_B) = \\ = {}^A \hat{x}_B^T ({}^A \hat{y}_B \times {}^A \hat{z}_B) \end{array} \right.$$

Special Orthogonal Group

$$A \in SO(3)$$

$$= 1$$

- **Special Orthogonal Group:** Space of Rotation Matrices in \mathbb{R}^n is defined as

$$\underbrace{SO(3)} \quad \underline{SO}(n) = \{R \in \mathbb{R}^{n \times n} : R^T R = I, \det(R) = 1\}$$

- $SO(n)$ is a *group*. We are primarily interested in $SO(3)$ and $SO(2)$, rotation groups of \mathbb{R}^3 and \mathbb{R}^2 , respectively.
- **Group** is a set G , together with an operation \bullet , satisfying the following group axioms:
 - **Closure:** $a \in G, b \in G \Rightarrow a \bullet b \in G$
 - **Associativity:** $(a \bullet b) \bullet c = a \bullet (b \bullet c), \forall a, b, c \in G$
 - **Identity element:** $\exists e \in G$ such that $e \bullet a = a$, for all $a \in G$.
 - **Inverse element:** For each $a \in G$, there is a $b \in G$ such that $a \bullet b = b \bullet a = e$, where e is the identity element.

Use of Rotation Matrix (1/2)

- Representing an orientation ${}^A R_B$ from definition

- Changing the reference frame ${}^A R_B$. Given vector v , it's coordinates in $\{A\}$, $\{B\}$, are ${}^A v$, ${}^B v$

${}^A v = {}^A R_B {}^B v$ \therefore "coordinate-free" proof

— same physical vector v , suppose ${}^A v = \begin{bmatrix} \alpha_1 \\ \alpha_2 \\ \alpha_3 \end{bmatrix}$, ${}^B v = \begin{bmatrix} \beta_1 \\ \beta_2 \\ \beta_3 \end{bmatrix}$

"physics"
 $\Rightarrow v = \alpha_1 \hat{x}_A + \alpha_2 \hat{y}_A + \alpha_3 \hat{z}_A$

$$v = \beta_1 \hat{x}_B + \beta_2 \hat{y}_B + \beta_3 \hat{z}_B$$

$$\Rightarrow \alpha_1 \hat{x}_A + \alpha_2 \hat{y}_A + \alpha_3 \hat{z}_A = \beta_1 \hat{x}_B + \beta_2 \hat{y}_B + \beta_3 \hat{z}_B \quad \text{— "physics"}$$

"state this physics" in $\{A\}$ -frame $\Rightarrow \alpha_1 \hat{x}_A + \alpha_2 \hat{y}_A + \alpha_3 \hat{z}_A = \beta_1 \hat{x}_B + \beta_2 \hat{y}_B + \beta_3 \hat{z}_B$

Use of Rotation Matrix (2/2)



$$\underbrace{\begin{bmatrix} {}^A\hat{x}_A & {}^A\hat{y}_A & {}^A\hat{z}_A \end{bmatrix}}_{\mathbf{I}} \cdot \underbrace{\begin{bmatrix} \alpha_1 \\ \alpha_2 \\ \alpha_3 \end{bmatrix}}_{\mathbf{A}_V} = \underbrace{\begin{bmatrix} {}^A\hat{x}_B & {}^A\hat{y}_B & {}^A\hat{z}_B \end{bmatrix}}_{\mathbf{A}_B} \underbrace{\begin{bmatrix} \beta_1 \\ \beta_2 \\ \beta_3 \end{bmatrix}}_{\mathbf{B}_V}$$

$$\mathbf{I} \cdot \mathbf{A}_V = \mathbf{A}_B \mathbf{B}_V$$

- Rotating a vector or a frame $\text{Rot}(\hat{\omega}, \theta)$: will be discussed in next lecture.

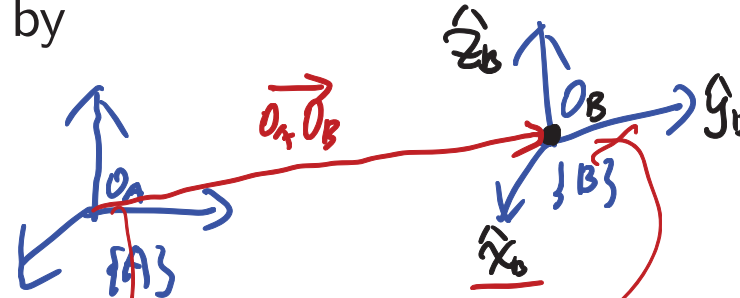
"action": verb operator view

Rigid Body Configuration

different orientation

- Given two coordinate frames {A} and {B}, the configuration of B relative to A is determined by

- ${}^A R_B$ and ${}^A O_B$



$${}^A R_B = \begin{bmatrix} {}^A \hat{x}_B & {}^A \hat{y}_B & {}^A \hat{z}_B \end{bmatrix}$$

${}^A O_B$

- For a (free) vector r , its coordinates ${}^A r$ and ${}^B r$ are related by:

r

$${}^A r = {}^A R_B {}^B r$$

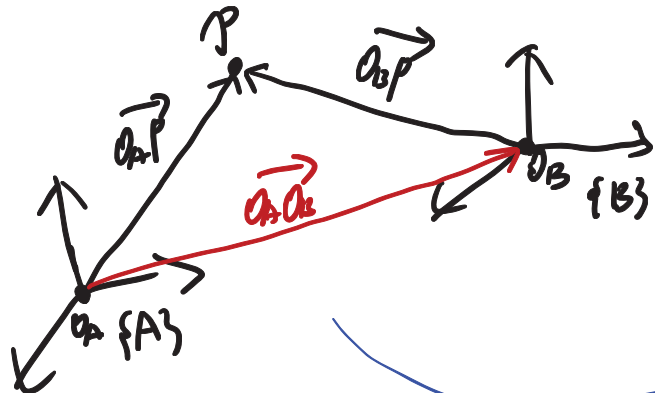
Linear relation

~~$${}^A \tilde{r} = \begin{bmatrix} {}^A r \\ 1 \end{bmatrix}, \quad {}^B \tilde{r} = \begin{bmatrix} {}^B r \\ 1 \end{bmatrix}$$

$${}^A \tilde{r} = {}^A T_B {}^B \tilde{r}$$

$$\Rightarrow {}^A \tilde{r} = {}^A O_B + {}^A R_B {}^B \tilde{r}$$~~

- For a point p , its coordinates ${}^A p$ and ${}^B p$ are related by:



"coordinate-free"

$$\vec{O_A P} = \vec{O_A O_B} + \vec{O_B P}$$

choose "B" frame to express "physics"

$${}^A p = {}^A O_B + {}^A R_B {}^B p$$

$${}^A \vec{O_A P} = {}^A \vec{O_A O_B} + {}^A \vec{O_B P} = {}^A R_B {}^B p$$

Homogeneous Transformation Matrix

- Homogeneous Transformation Matrix: ${}^A T_B$

$${}^A p = \underbrace{{}^A O_a}_{\in \mathbb{R}^3} + \underbrace{{}^A R_B}_{3 \times 3} \cdot \underbrace{{}^B p}_{\in \mathbb{R}^3} \iff \begin{bmatrix} {}^A p \\ 1 \end{bmatrix} = \underbrace{\begin{bmatrix} {}^A R_B & | & {}^A O_B \\ \hline 0 & | & 1 \end{bmatrix}}_{4 \times 4} \begin{bmatrix} {}^B p \\ 1 \end{bmatrix}$$

$\swarrow \mathbb{R}^3$
 $\underbrace{\quad}_{1 \times 3}$ $\underbrace{\quad}_{1 \times 1}$

affine relation

$${}^A T_B \triangleq \begin{bmatrix} {}^A R_B & | & {}^A O_B \\ \hline 0 & | & 1 \end{bmatrix} \quad \left(T = (R, p), \text{ configuration of } \{B\} \text{ relative to } \{A\} \right)$$

- Homogeneous coordinates:

Given a point $p \in \mathbb{R}^3$, its homogeneous coordinate is given

by $\tilde{p} = \begin{bmatrix} p \\ 1 \end{bmatrix} \in \mathbb{R}^4 \Rightarrow \boxed{{}^A \tilde{p} = {}^A T_B \cdot {}^B \tilde{p}}$

Given vector v , its homogeneous coordinate is $\tilde{v} = \begin{bmatrix} v \\ 0 \end{bmatrix}$

$$v = p_1 - p_2 \quad \tilde{v} = \tilde{p}_1 - \tilde{p}_2$$

Example of Homogeneous Transformation Matrix

Fixed frame $\{a\}$; end effector frame $\{b\}$, the camera frame $\{c\}$, and the workpiece frame $\{d\}$. Suppose $\|p_c - p_b\| = 4$

1. Camera "location"? ${}^aR_c, {}^ap_c$

$${}^aT_c = ({}^aR_c, {}^ap_c)$$

$${}^aR_c = \begin{bmatrix} 0 & 1 & 0 \\ 1 & 0 & 0 \\ 0 & 0 & -1 \end{bmatrix}, \quad {}^ap_c = \begin{bmatrix} -1 \\ 1 \\ 2 \end{bmatrix}$$

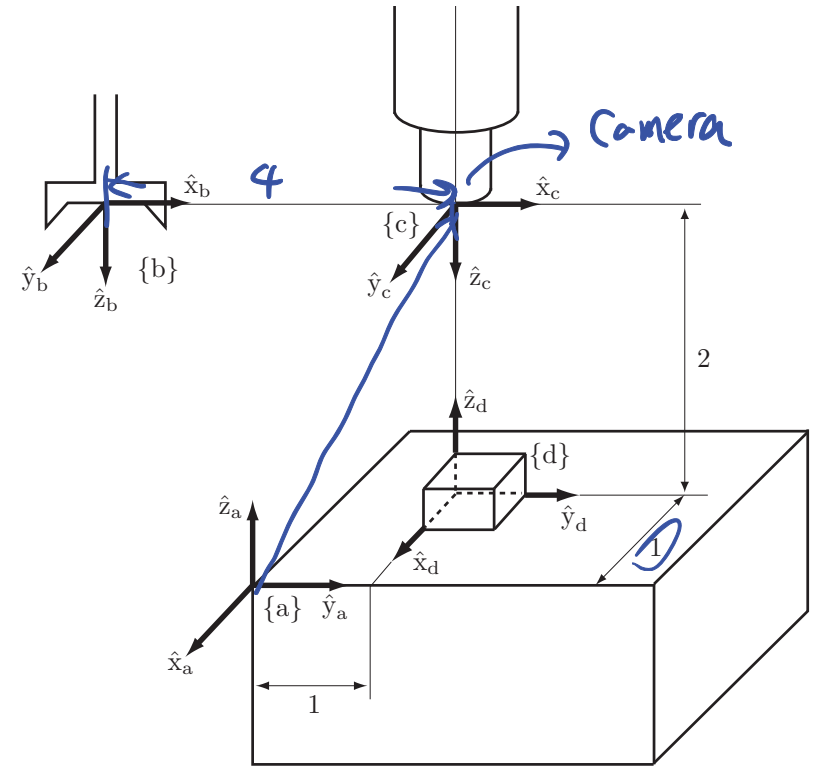
$$\uparrow \quad \boxed{\hat{x} = 0 \cdot \hat{x}_a + 1 \cdot \hat{y}_a + 0 \cdot \hat{z}_a}$$

$${}^aT_c = \left[\begin{array}{ccc|c} {}^aR_c & & & {}^ap_c \\ \hline 0 & 0 & 0 & 1 \end{array} \right]$$

2. end-effector frame: aT_b

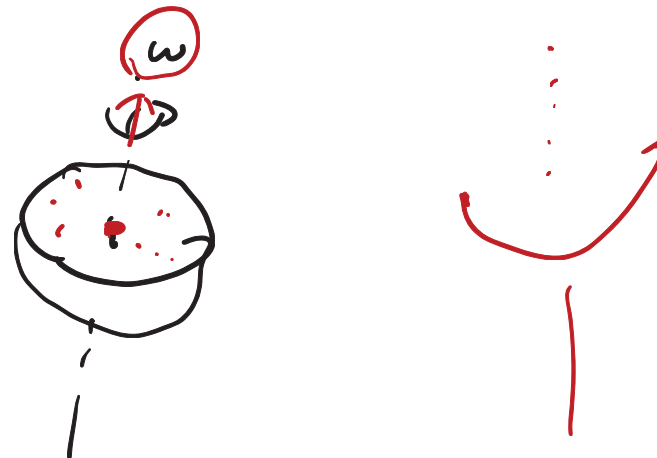
$${}^aT_b = ({}^aT_c) {}^cT_b$$

$${}^cT_b = \begin{bmatrix} 1 & 0 & 0 & -4 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ \hline 0 & 0 & 0 & 1 \end{bmatrix}$$



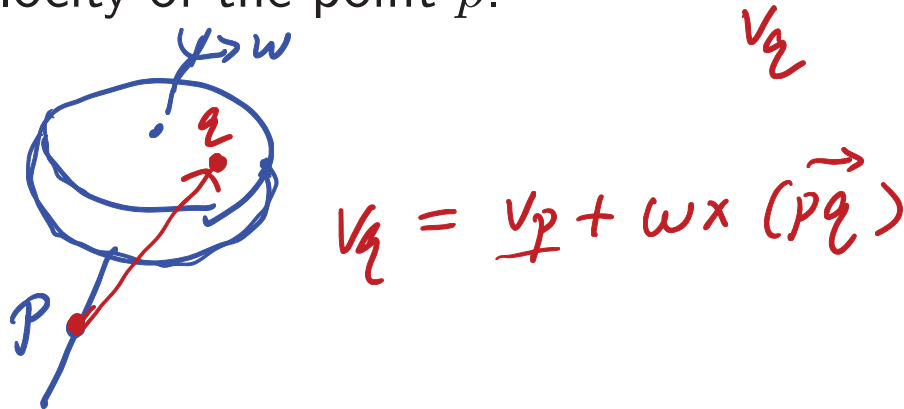
Outline

- Rigid Body Configuration
- Rigid Body Velocity (Twist) ←
- Geometric Aspect of Twist: Screw Motion



Rigid Body Velocity (1/2)

- Consider a rigid body with angular velocity: ω (this is a free vector).
- Suppose the actual rotation axis passes through a point p ; Let v_p be the velocity of the point p .



Question: A rigid body contains infinitely many points with different velocities. How to parameterize all of their velocities?

- Consider an arbitrary body-fixed point q (means that the point is rigidly attached to the body, and moves with the body), we have:

$$\boxed{v_q = v_p + \omega \times (\vec{pq})} \quad (3)$$

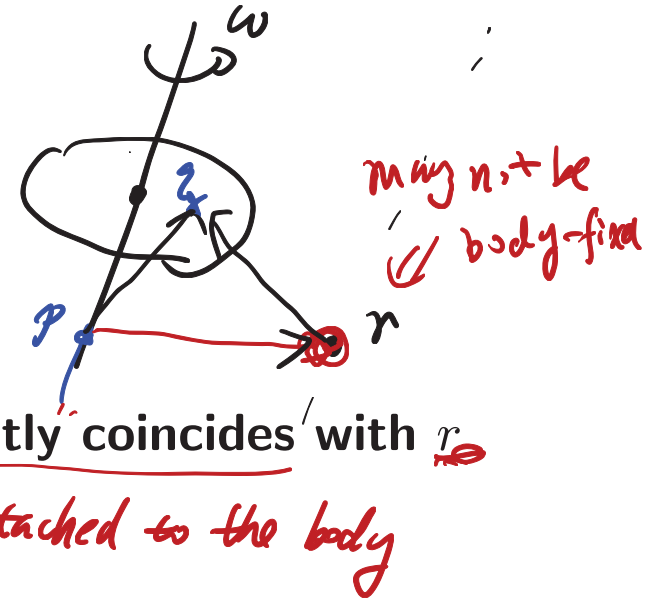
- The velocity of an arbitrary body-fixed point depends only on $(\underline{\omega}, \underline{v}_p, p)$ and the location of the point q .

Rigid Body Velocity (2/2)

- Fact: The representation form (3) is independent of the reference point p .

- Consider an arbitrary point r in space

- r may not be on the rotation axis
- r may be a stationary point in space (does not move)



- Let v_r be the velocity of the body-fixed point currently coincides with r

- We still have: $v_q = v_r + \omega \times (\vec{r}q)$

- $$v_q = v_p + \omega \times (\vec{p}q)$$

$$= v_r - \omega \times \vec{pr} + \omega \times (\vec{p}q)$$

$$= v_r + \omega \times (\vec{p}q - \vec{pr}) = v_r + \omega \times \vec{rq}$$

- The body can be regarded as translating with a linear velocity v_r , while rotating with angular velocity ω about an axis passing through r

Rigid Body Velocity: Spatial Velocity (Twist)

Featherstone

Murray

- Spatial Velocity (Twist): $\mathcal{V}_r = (\omega, v_r) \in \mathbb{R}^6$
 - ω angular velocity
 - v_r velocity of the body-fixed point currently coincides with r
 - For any other body-fixed point q , its velocity is

$$v_q = v_r + \omega \times (\vec{r}_q)$$

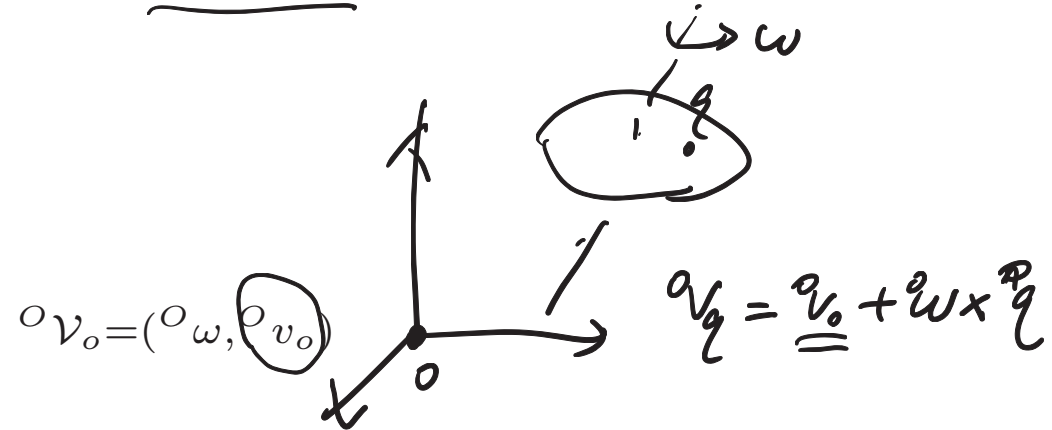
- Twist is a “physical” quantity (just like linear or angular velocity)
 - It can be represented in any frame for any chosen reference point r
- A rigid body with $\mathcal{V}_r = (\omega, v_r)$ can be “thought of” as translating at v_r while rotating with angular velocity ω about an axis passing through r
 - This is just one way to interpret the motion.

Spatial Velocity Representation in a Reference Frame

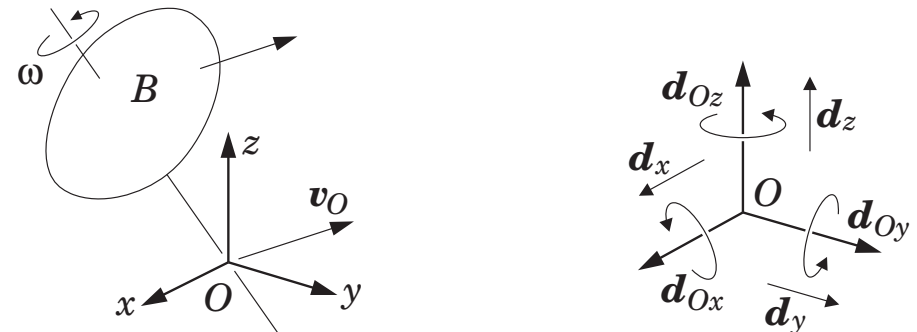
- Given frame $\{O\}$ and a spatial velocity $\underline{\mathcal{V}}$
- Choose o (the origin of $\{O\}$) as the reference point to represent the rigid body velocity

- Coordinates for the \mathcal{V} in $\{O\}$:

$A_{\mathcal{V}}$



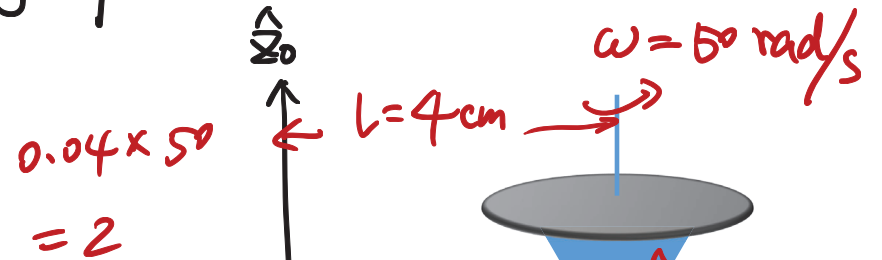
- By default, we assume the origin of the frame is used as the reference point:
 ${}^o\mathcal{V} = {}^o\mathcal{V}_o$



Example of Twist I

- Example I: What's the twist of the spinning top?

choose $\{0\}$ -frame. ${}^0v_{top} = \begin{bmatrix} {}^0\omega \\ \vdots \\ {}^0v_b \\ \vdots \\ 0 \\ 0 \\ 0 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 50 \text{ rad/s} \\ \vdots \\ 0 \\ 0 \\ 0 \end{bmatrix}$



choose $\{A\}$ -frame: ${}^A v = \begin{bmatrix} {}^A \omega \\ \vdots \\ {}^A v_{OA} \\ \vdots \\ 0 \\ 0 \\ 0 \end{bmatrix}$

$${}^0T_A = \begin{bmatrix} {}^0R_A & {}^0p_A \\ 0 & 1 \end{bmatrix} \Rightarrow {}^0X_A = \begin{bmatrix} R & 0 \\ (p)R & R \end{bmatrix}$$

V.F.T.

$${}^0v_{top} \in \mathbb{R}^6 = \underbrace{{}^0X_A}_{6 \times 6} \underbrace{{}^A v_{top}}_{\in \mathbb{R}^6}$$

Example of Twist II

- Example II: what's

$${}^b v_{car} \quad {}^s v_{car}$$

Car rotates about \hat{w}

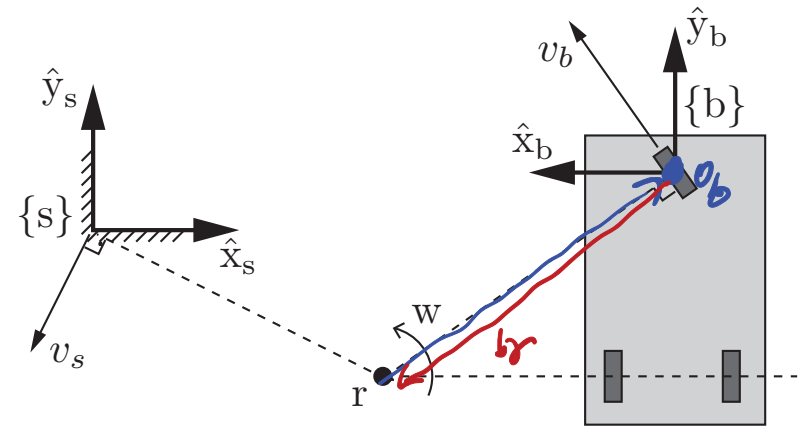
$${}^b v_{car} = \begin{bmatrix} {}^b \omega \\ {}^b v_o \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ -2 \\ \hline 2.8 \\ 4 \\ 0 \end{bmatrix}$$

$${}^b \omega = \begin{bmatrix} 0 \\ 0 \\ -2 \end{bmatrix}$$

$${}^b v_o = {}^b \omega \times \underbrace{({}^b r_o)}_{=(-r)} = \begin{bmatrix} 0 \\ 0 \\ -2 \end{bmatrix} \times \begin{bmatrix} -2 \\ 4 \\ 0 \end{bmatrix} = \begin{bmatrix} 2.8 \\ 4 \\ 0 \end{bmatrix}$$

$${}^s v_{car} = \begin{bmatrix} {}^s \omega \\ {}^s v_o \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ -2 \\ \hline -2 \\ -4 \\ 0 \end{bmatrix}$$

$$\begin{aligned} {}^s v_o &= {}^s \omega \times ({}^s r_s) \\ &= {}^s \omega \times (-r) \\ &= \begin{bmatrix} -2 \\ -4 \\ 0 \end{bmatrix} \end{aligned}$$



$$r_s = (2, -1, 0), \quad r_b = (2, -1.4, 0), \quad w = 2 \text{ rad/s}$$

Change Reference Frame for Twist (1/2)

- Given a twist \mathcal{V} , let ${}^A\mathcal{V}$ and ${}^B\mathcal{V}$ be their coordinates in frames $\{A\}$ and $\{B\}$

$${}^A\mathcal{V} = \begin{bmatrix} \underline{{}^A\omega} \\ \underline{{}^A v_A} \end{bmatrix}, \quad \textcircled{{}^B\mathcal{V}} = \begin{bmatrix} \underline{{}^B\omega} \\ \underline{{}^B v_B} \end{bmatrix}$$

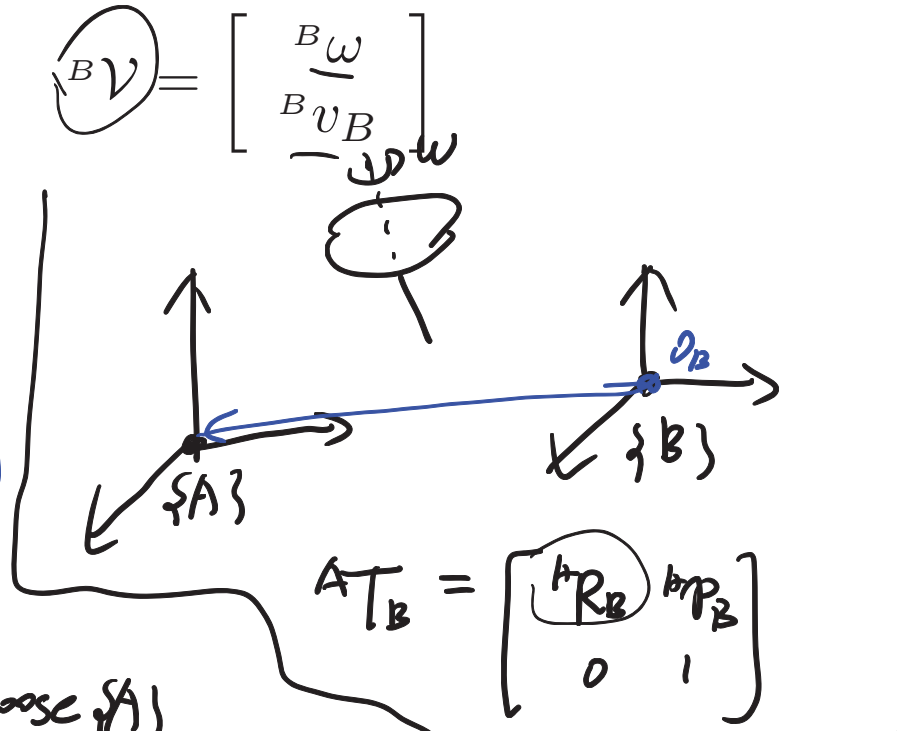
- They are related by ${}^A\mathcal{V} = {}^A X_B {}^B\mathcal{V}$

① $\underline{{}^A\omega} = {}^A R_B \underline{{}^B\omega} \implies \underline{{}^A\omega} = \begin{bmatrix} {}^A R_B & 0 \\ 0 & 1 \end{bmatrix} \begin{bmatrix} \underline{{}^B\omega} \\ \underline{{}^B v_B} \end{bmatrix}$

② "coordinate-free"
 v_A : body-fixed pt coincides with O_A
 velocity

v_B : O_B

$$v_A = v_B + \omega \times \underline{{}^A(BA)}$$



\implies choose $\{A\}$ to express "physics"

$${}^A v_A = \underline{{}^A v_B} + \underline{{}^A \omega} \times \underline{{}^A(BA)} \quad \xrightarrow{\quad} \quad = \underline{{}^A(-O_B)}$$

$$= {}^A R_B \underline{{}^B v_B} + \underline{{}^A R_B \omega} \times \underline{{}^A(-O_B)}$$

$$= {}^A R_B \underline{{}^B v_B} + \underline{{}^A O_B} \times ({}^A R_B \underline{{}^B \omega})$$

\Downarrow $[{}^A O_B]$

Change Reference Frame for Twist (2/2)

$$\Rightarrow A_j = \begin{bmatrix} A_W \\ A_V_A \end{bmatrix} = \begin{bmatrix} {}^A R_B & 0 \\ \underbrace{[{}^A O_B] {}^A R_B}_{6 \times 6 \text{ matrix}} & {}^A R_B \end{bmatrix} \begin{bmatrix} B_W \\ B_V_B \end{bmatrix} = \begin{bmatrix} [{}^A O_B] {}^A R_B & {}^A R_B \end{bmatrix} \begin{bmatrix} B_W \\ B_V_B \end{bmatrix}$$

$\triangleq A X_B$

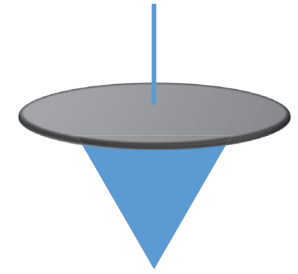
${}^A T_B \triangleq (R, p) \Rightarrow A X_B = \begin{bmatrix} R & 0 \\ [p]R & R \end{bmatrix}$

- If configuration $\{B\}$ in $\{A\}$ is $T = (R, p)$, then

$${}^A X_B = [Ad_T] \triangleq \begin{bmatrix} R & 0 \\ [p]R & R \end{bmatrix}$$

Given $T = (R, p) \xrightarrow{[Ad_T]} X =$

Example I Revisited



Outline

- Rigid Body Configuration
- Rigid Body Velocity (Twist)

- Geometric Aspect of Twist: Screw Motion

- recall: • linear velocity $v = \begin{bmatrix} 1 \\ 2 \end{bmatrix} = \underbrace{\|v\|}_{\text{scalar}} \cdot \hat{v} = \sqrt{5} \cdot \begin{bmatrix} \frac{1}{\sqrt{5}} \\ \frac{2}{\sqrt{5}} \end{bmatrix}$

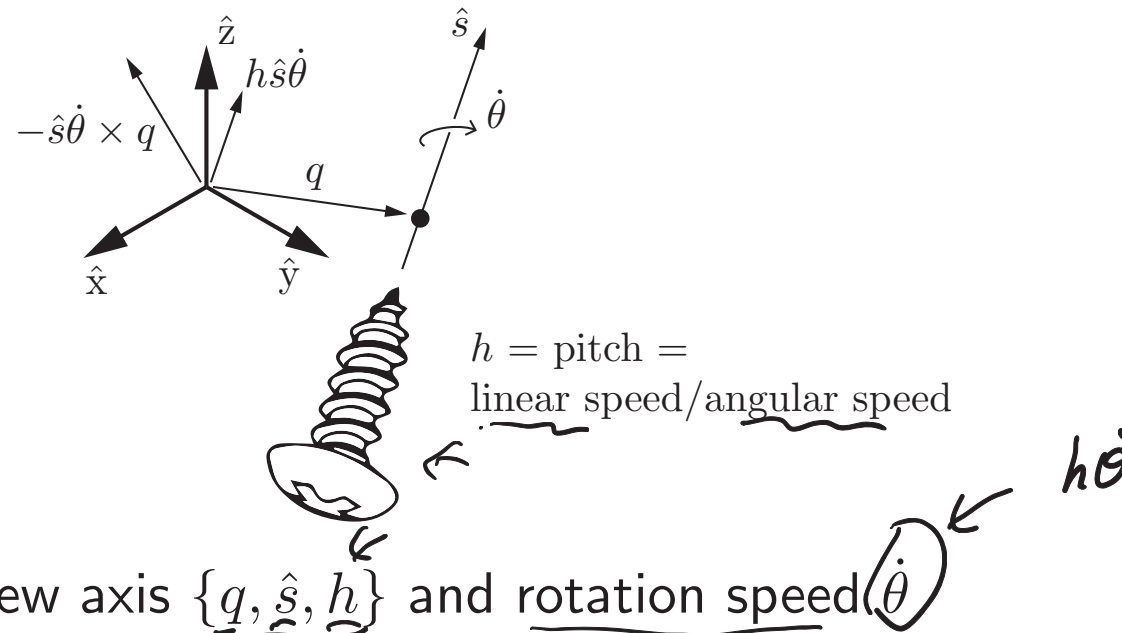
- angular velocity: $\omega = \underbrace{\hat{\omega}}_{\substack{\downarrow \\ \text{rotation axis}}} \cdot \underbrace{\dot{\theta}}_{\substack{\uparrow \\ \text{scalar}}}$ (unit vector)

• rigid-body velocity: twist $\mathcal{V} = (v, \omega)$

Screw Motion: Definition

→ : standard/canonical motion for rigid body motion

- Rotating about an axis while also translating along the axis



- Represented by screw axis $\{q, \hat{s}, h\}$ and rotation speed $\dot{\theta}$
 - \hat{s} : unit vector in the direction of the rotation axis
 - q : any point on the rotation axis
 - h : **screw pitch** which defines the ratio of the linear velocity along the screw axis to the angular velocity about the screw axis
- Theorem (Chasles): Every rigid body motion can be realized by a screw motion.

From Screw Motion to Twist

- Consider a rigid body under a screw motion with screw axis $\{\hat{s}, h, q\}$ and (rotation) speed $\dot{\theta}$

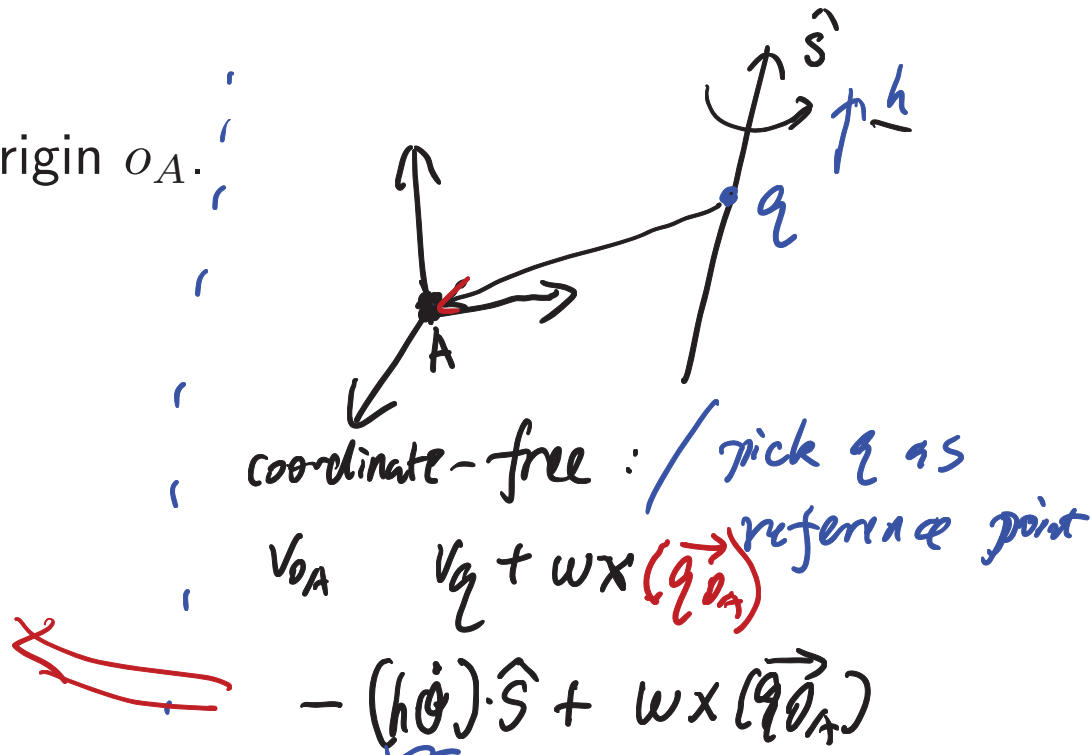
- Fix a reference frame $\{A\}$ with origin o_A .

- Find the twist ${}^A\mathcal{V} = ({}^A\omega, {}^A v_{o_A})$

$${}^A\omega = {}^A\hat{s} \cdot \dot{\theta}$$

$${}^A v_{o_A} = (h\dot{\theta}) \hat{s} + {}^A\omega \times (-{}^A q)$$

$$= {}^A\hat{s} (h\dot{\theta}) - {}^A\omega \times {}^A q$$



- Result:** given screw axis $\{\hat{s}, h, q\}$ with rotation speed $\dot{\theta}$, the corresponding twist $\mathcal{V} = (\omega, v)$ is given by

$$\omega = \hat{s} \dot{\theta} \quad v = -\hat{s} \dot{\theta} \times q + h \hat{s} \dot{\theta}$$

- The result holds as long as all the vectors and the twist are represented in the same reference frame

From Twist to Screw Motion

- The converse is true as well: given any twist $\mathcal{V} = (\omega, v)$ we can always find the corresponding screw motion $\{q, \hat{s}, h\}$ and $\dot{\theta}$

- If $\omega = 0$, then it is a pure translation ($h = \infty$)

pure linear

$$\hat{s} = \frac{v}{\|v\|}, \quad \dot{\theta} = \|v\|, \quad h = \infty, \quad q \text{ can be arbitrary}$$

- If $\omega \neq 0$:

$$\hat{s} = \frac{\omega}{\|\omega\|}, \quad \dot{\theta} = \|\omega\|, \quad q = \frac{\omega \times v}{\|\omega\|^2}, \quad h = \frac{\omega^T v}{\|\omega\|}$$

you can plug into the eqn on previous slide to verify the result.

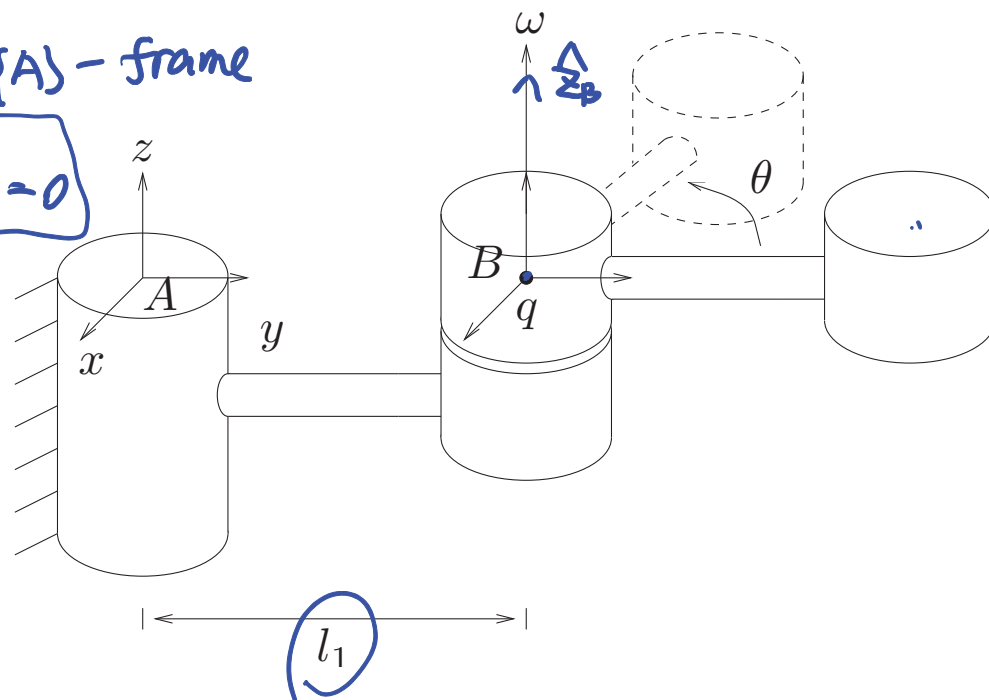
Examples: Screw Axis and Twist

- What is the twist that corresponds to rotating about \hat{z}_B with $\dot{\theta} = 2$? *choose {A} - frame*

screw axis: $A \hat{s} = A \hat{z}_B = \begin{bmatrix} 0 \\ 0 \\ 1 \end{bmatrix}, A q = \begin{bmatrix} 0 \\ l_1 \\ 0 \end{bmatrix}, h = 0$
 $\dot{\theta} = 2$

$${}^A \omega = {}^A \hat{s} \dot{\theta} = \begin{bmatrix} 0 \\ 2 \\ 0 \end{bmatrix}$$

$${}^A v_{O_A} = 0 - {}^A \omega \times {}^A q = - \begin{bmatrix} 0 \\ 2 \\ 0 \end{bmatrix} \times \begin{bmatrix} 0 \\ l_1 \\ 0 \end{bmatrix} = \begin{bmatrix} 2l_1 \\ 0 \\ 0 \end{bmatrix}$$



- What is the screw axis for twist $\mathcal{V} = (0, 2, 2, 4, 0, 0)$?

$$\omega = \begin{bmatrix} 0 \\ 2 \\ 2 \end{bmatrix}, v = \begin{bmatrix} 4 \\ 0 \\ 0 \end{bmatrix}, \text{ use the previous formula}$$

$$\Rightarrow q = \frac{\omega \times v}{\|\omega\|^2} = \#$$

Screw Representation of a Twist

$$\omega = \hat{\omega} \dot{\theta}$$

- Recall: an angular velocity vector ω can be viewed as $\hat{\omega} \dot{\theta}$, where $\hat{\omega}$ is the unit rotation axis and $\dot{\theta}$ is the rate of rotation about that axis
- Similarly, a twist (spatial velocity) \mathcal{V} can be interpreted in terms of a screw axis \hat{S} and a velocity $\dot{\theta}$ about the screw axis
- Consider a rigid body motion along a screw axis $\hat{S} = \{\hat{s}, h, q\}$ with speed $\dot{\theta}$. With slight abuse of notation, we will often write its twist as

$$\mathcal{V} = \hat{S} \dot{\theta}$$

← spatial velocity (ω, v)

- In this notation, we think of \hat{S} as the twist associated with a unit speed motion along the screw axis $\{\hat{s}, h, q\}$

More Discussions